

Fundamentals of Simulation Methods

Exercise Sheet 5

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Multigrid solver for linear systems

1 Galerkin coarse grid approximation

Consider the one-dimensional problem

$$\frac{\partial^2 \Phi}{\partial^2 x} = 4\pi G \rho(x), \quad (1)$$

which we want to solve using a multigrid accelerated iterative method on a grid of size $N = 8$ with spacing h . The problem can be rephrased in the form

$$\mathbf{A}\mathbf{x} = \mathbf{b}, \quad (2)$$

where $\mathbf{x} \hat{=} \Phi$ and $\mathbf{b} \hat{=} 4\pi G \rho h^2$ are vectors of size N and \mathbf{A} is an $N \times N$ -matrix. Given the operator

$$\mathbf{A}^{(h)} = \begin{pmatrix} -2 & 1 & 0 & 0 & 0 & 0 & 0 & 1 \\ 1 & -2 & 1 & 0 & 0 & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & 0 & 1 & -2 & 1 \\ 1 & 0 & 0 & 0 & 0 & 0 & 1 & -2 \end{pmatrix} \quad (3)$$

on the finest grid level ($N = 8$), find the operator $\mathbf{A}^{(2h)}$ on the next coarser grid level ($N = 4$, spacing $2h$). In order to do so, carry out the following steps:

(a) Find the prolongation and restriction operators in matrix form,

$$\mathbf{I}_{(2h)}^{(h)} : \begin{cases} x_{2i}^{(h)} = x_i^{(2h)}, \\ x_{2i+1}^{(h)} = \frac{1}{2} (x_i^{(2h)} + x_{i+1}^{(2h)}), \end{cases} \quad (4)$$

$$\mathbf{I}_{(h)}^{(2h)} : x_i^{(2h)} = \frac{1}{4} (x_{2i-1}^{(h)} + 2x_{2i}^{(h)} + x_{2i+1}^{(h)}), \quad (5)$$

where $0 \leq i < N/2$. Hint: $\mathbf{I}_{(2h)}^{(h)} \mathbf{x}^{(h)} = \mathbf{x}^{(2h)}$ and vice versa.

(b) Now calculate the operator $\mathbf{A}^{(2h)}$ on the coarse grid using the Galerkin coarse grid approximation $\mathbf{A}_{\text{Gal}}^{(2h)} = \mathbf{I}_{(h)}^{(2h)} \mathbf{A}^{(h)} \mathbf{I}_{(2h)}^{(h)}$.

(c) Compare the result to an operator obtained by direct discretization on the coarse grid.

- (a) The notation employed in part (a) is quite confusing. We attempt a more natural way of expressing eqs. (4) and (5) by introducing

$$\begin{aligned}\mathbf{I}_{(2h)}^{(h)} &\equiv \mathbf{P}^{(2h \rightarrow h)}, \\ \mathbf{I}_{(h)}^{(2h)} &\equiv \mathbf{R}^{(h \rightarrow 2h)},\end{aligned}\tag{6}$$

and resorting to the well-known index notation:

$$x_i^{(h)} = \sum_{j=1}^N P_{ij}^{(2h \rightarrow h)} x_j^{(2h)},\tag{7}$$

$$x_i^{(2h)} = \sum_{j=1}^N R_{ij}^{(h \rightarrow 2h)} x_j^{(h)}.\tag{8}$$

Looking at eqs. (7) and (8), it is now clear that $\mathbf{P}^{(2h \rightarrow h)}$ and $\mathbf{R}^{(h \rightarrow 2h)}$ are $N \times N/2$ - and $N/2 \times N$ -matrices, respectively.

To implement eqs. (4) and (5), $\mathbf{P}^{(2h \rightarrow h)}$ and $\mathbf{R}^{(h \rightarrow 2h)}$ must take the form

$$\mathbf{P}^{(2h \rightarrow h)} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 1/2 & 1/2 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 1/2 & 1/2 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 1/2 & 1/2 \\ 0 & 0 & 0 & 1 \\ 1/2 & 0 & 0 & 1/2 \end{pmatrix},\tag{9}$$

$$\mathbf{R}^{(h \rightarrow 2h)} = \begin{pmatrix} 1/2 & 1/4 & 0 & 0 & 0 & 0 & 0 & 1/4 \\ 0 & 1/4 & 1/2 & 1/4 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1/4 & 1/2 & 1/4 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1/4 & 1/2 & 1/4 \end{pmatrix}\tag{10}$$

- (b) Multiplying the operator $\mathbf{A}^{(h)}$ from the left with a restriction and from the right with a prolongation, we find

$$\mathbf{A}_{\text{Gal}}^{(2h)} = \mathbf{R}^{(h \rightarrow 2h)} \mathbf{A}^{(h)} \mathbf{P}^{(2h \rightarrow h)} = \begin{pmatrix} -1/2 & 1/4 & 0 & 1/4 \\ 1/4 & -1/2 & 1/4 & 0 \\ 0 & 1/4 & -1/2 & 1/4 \\ 1/4 & 0 & 1/4 & -1/2 \end{pmatrix} = \frac{1}{4} \begin{pmatrix} -2 & 1 & 0 & 1 \\ 1 & -2 & 1 & 0 \\ 0 & 1 & -2 & 1 \\ 1 & 0 & 1 & -2 \end{pmatrix}.\tag{11}$$

- (c) Setting out from the coarse grid instead with only $N/2$ points per dimension spaced $2h$ apart, the discretized Poisson equation would be given by

$$\frac{\Phi_{i+1} - 2\Phi_i + \Phi_{i-1}}{(2h)^2} = 4\pi G\rho_i,\tag{12}$$

where the index i now runs from $i \in \{0, 1, 2, \dots, \frac{N}{2} - 1\}$ if N is still taken to be the number of points on the fine grid. If we then write this Poisson equation in the form of eq. (2), i.e. as a system of linear equations with the components of \mathbf{b} now given by

$$b_i = 4\pi G(2h^2)\rho_i,\tag{13}$$

we see that $\mathbf{A}^{(2h)}$ has to be

$$\mathbf{A}^{(2h)} = \begin{pmatrix} -2 & 1 & 0 & 1 \\ 1 & -2 & 1 & 0 \\ 0 & 1 & -2 & 1 \\ 1 & 0 & 1 & -2 \end{pmatrix} = 4\mathbf{A}_{\text{Gal}}^{(2h)}. \quad (14)$$

This $\mathbf{A}^{(2h)}$ is larger by a factor of 4 than the one we found in part (b), as it should be since its corresponding \mathbf{b} is also four times larger.

2 Iteratively solving Poisson's equation

- (a) Suppose we want to solve for the gravitational potential Φ on a *periodic* mesh with square-shaped domain $[0, L] \times [0, L]$ (with $L = 1$ for definiteness) containing $N \times N$ cells and housing the mass density distribution

$$\rho(\mathbf{r}) = \rho_0 \exp\left(-\frac{r^2}{2\eta^2}\right), \quad (15)$$

where $\eta = \frac{L}{10}$ is the spread of the mass spike, and $\rho_0 = 10$. Discretization of the Poisson equation gives an iteration rule where the left-hand side gives the updated values in terms of the old values on the right-hand side, in the form

$$x_{i,j}^{(n+1)} = \frac{1}{4} \left(x_{i-1,j}^{(n)} + x_{i+1,j}^{(n)} + x_{i,j-1}^{(n)} + x_{i,j+1}^{(n)} - b_{i,j} \right), \quad (16)$$

where $x_{i,j} = \Phi_{i,j}$ and $b_{i,j} = 4\pi Gh^2 \rho_{i,j} - \bar{b}$. This can be readily used for Jacobi or Gauss-Seidel iteration schemes. Write a function `jacobi_step(x[], b[], N)` that replaces the input array `x[]` with a correspondingly updated array after one iteration step.

- (b) To characterize the current error in the solution, write a function that calculates the residual

$$r_{i,j} = b_{i,j} - (\mathbf{A}\mathbf{x})_{i,j}. \quad (17)$$

Also, write a function that computes the norm

$$S = \left(\sum_{i,j} r_{i,j}^2 \right)^{1/2} \quad (18)$$

of this vector and returns it. We can use S as a simple measure of the overall error.

- (c) Now specialize to the case of $N = 256$ and write a routine that suitably initializes the density field, with the density peak placed in the center of the box. Set the initial guess for the potential to $\mathbf{x} = \mathbf{0}$. Carefully think about how to initialize the array \mathbf{b} (for a periodic system, eq. (2) is only solvable for $\langle \mathbf{b} \rangle = \mathbf{0}$ because the matrix \mathbf{A} is singular, therefore \bar{b} has to be chosen appropriately). Now add a loop that calls the Jacobi iteration $N_{\text{steps}} = 2000$ times, and after each step, determines the norm of the residual. You may use the C or Python template on the Moodle site and fill in the missing parts in order to reduce the coding work, or write everything from scratch if you prefer. Make a plot of the decay of the log of this residual as a function of step number. What do you expect to get for $\sum_{i,j} \Phi_{i,j}$?
- (d) Now produce a second version of your program in which the Jacobi iteration is replaced by Gauss-Seidel iteration where new values for elements of \mathbf{x} are used as soon as they become available, replacing any old value in the array.

- (e) Change the Gauss-Seidel scheme to a red-black Gauss-Seidel iteration, i.e. you first update the red cells in a chess-board pattern overlaid over the mesh, then the black cells in a second pass. Produce a common plot of the decay rates of the residual for N_{steps} steps with the three iteration variants considered thus far. Interpret the results. The red-black scheme result will perhaps disappoint you at first – what’s going on?
- (f) We now solve the problem with multigrid acceleration. First, write a function that carries out a restriction step of a mesh with dimension $N \times N$ (where N is a power of 2) onto a coarser mesh with dimension $N/2 \times N/2$. Make each point in the coarser mesh a weighted average of neighboring points, according to the stencil

$$\mathbf{T} = \begin{bmatrix} 1/16 & 1/8 & 1/16 \\ 1/8 & 1/4 & 1/8 \\ 1/16 & 1/8 & 1/16 \end{bmatrix}. \quad (19)$$

Give this function a calling signature `do_restrict(N, fine[], NN, coarse[])` (only $NN = N/2$ is allowed when calling this function).

- (g) Next, write a function that carries out a prolongation step of a mesh of dimension $N \times N$ onto a finer mesh of dimension $2N \times 2N$. Each point in the coarser mesh is additively injected into the finer mesh with weights

$$\mathbf{U} = \begin{bmatrix} 1/4 & 1/2 & 1/4 \\ 1/2 & 1 & 1/2 \\ 1/4 & 1/2 & 1/4 \end{bmatrix}. \quad (20)$$

Note that points of the fine mesh that see a coarse point with weight $1/4$ will get contributions from 4 coarse points in total, similarly for the points with weight $1/2$, while the fine-mesh point with weight 1 simply inherits the value of one of the coarse mesh points. Give this function a calling signature `do_prolong(NN, coarse[], N, fine[])` (again, only $NN = N/2$ is allowed when calling this function).

- (h) Now write a function `do_v_cycle(x[], b[], N)` that carries out a V-cycle multigrid iteration on the current solution vector. The steps of the function should be
1. Do a Gauss-Seidel step.
 2. Calculate the residuum.
 3. Restrict the residuum to a coarser mesh $N' = N/2$, and scale it by a factor of 4 to effectively take care of the $1/h \leftarrow 1/(2h)$ scaling of the differential operator (whose explicit form we leave invariant, for simplicity).
 4. Call `do_v_cycle` recursively for the coarser mesh, with a zero error vector as starting guess for \mathbf{x} , and the coarsened residual for \mathbf{b} .
 5. Now prolong the returned error vector to the fine mesh N .
 6. Add this to the current solution vector on the fine mesh.
 7. Do another Gauss-Seidel step.

Steps 2 to 6 are only done provided $N > 4$. To simplify the coding, you may use one of the code templates provided on the Moodle site.

- (i) Finally, solve the original problem by repeatedly calling your V-cycle iteration. Again, plot the residual as a function of the number of steps taken, and compare with the results obtained earlier for plain Jacobi and Gauss-Seidel iteration.

For parts (a) and (b), see `jacobi.c`.

(c) A half-logarithmic plot of the decay rate of the norm S of the residual is shown in fig. 1.

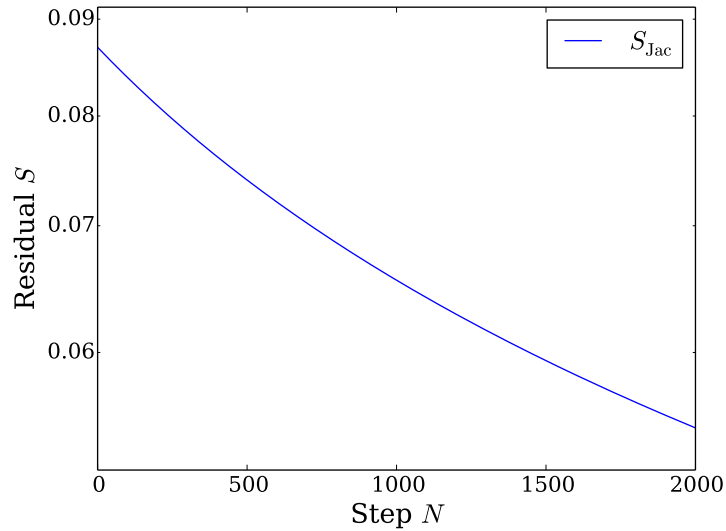


Figure 1: Decay of Jacobian residual norm S

(d) See `gausseidel.c`.

(e) Figure 2 displays the decay rates of the residual norm as produced by the Jacobi method and the Gauss-Seidel method with and without chessboard scheme.

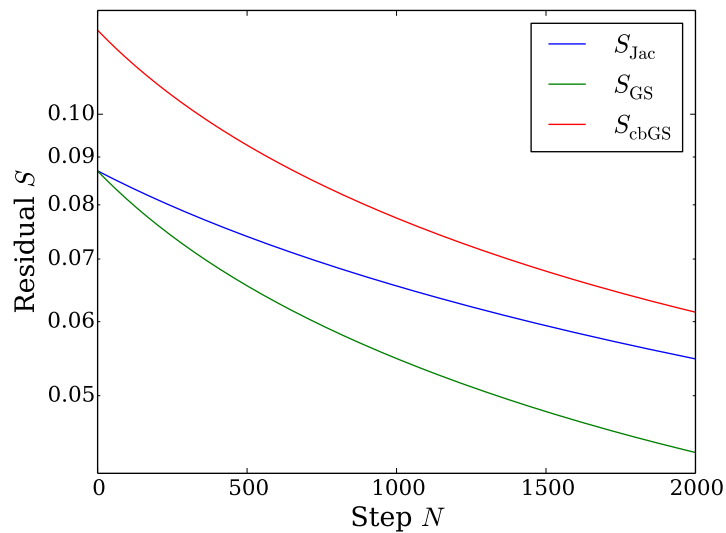


Figure 2: Gauss-Seidel residual decay rates compared to to that of the Jacobi method

As expected, we see that Gauss-Seidel's superior utilization of available information leads to the norm of its residual S_{GS} decaying considerably faster than that of Jacobi. The same improvement is not observed by further implementing a chessboard-like update scheme for the potential. The decay rate appears largely unchanged except for a spike during the very first step suggesting an ill-suited initial density distribution for this method.

For parts (f) to (h), see `multigrid.c`.

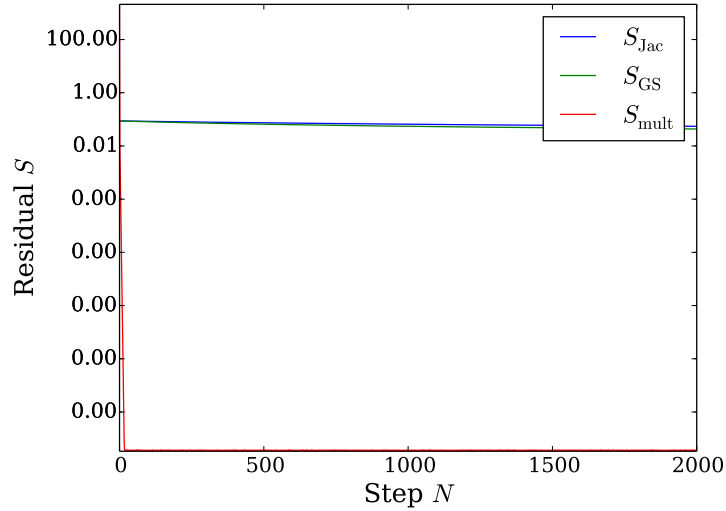


Figure 3: Residual decay of the multigrid method compared to Jacobi and Gauss-Seidel

- (i) Figure 3 compares the decay rate produced by the multigrid method to those of the Jacobi and Gauss-Seidel methods. We indeed find a dramatic increase in the convergence rate.

Note: A surface plot of the initial density distribution $\rho_i(\mathbf{r})$ over the entire simulation domain is shown in fig. 4a. Figures 4b to 4d display the noticeably differing corresponding potentials according to the Jacobi, Gauss-Seidel and multigrid methods, respectively. Since each method converges at a different rate, plotting the potential after a constant number of 2000 sweeps is akin to plotting it at different times.

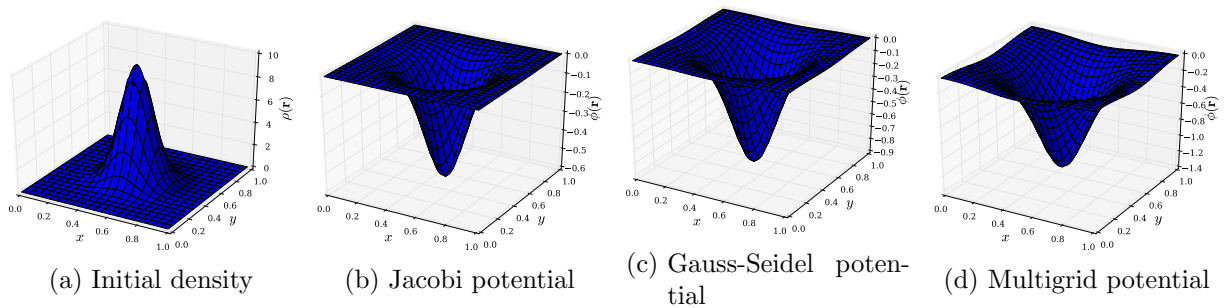


Figure 4: Surface plots of the initial density and calculated potentials